

Holographic Josephson Junction in 3+1 dimensions

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ABSTRACT: In arXiv:1101.3326[hep-th], a (2+1)-dimensional holographic Josephson junction was constructed, and it was shown that the DC Josephson current is proportional to the sine of the phase difference across the junction. In this paper, we extend this study to a holographic description for the (3+1)-dimensional holographic DC Josephson junction. By solving numerically the coupled differential equations, we also obtain the familiar characteristics of Josephson junctions.

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1 Introduction

The study of the AdS/CFT correspondence and its applications have been the subject of great activity since the original conjecture was made by Maldacena [1]. The AdS/CFT correspondence is a powerful tool to study the strong coupled field theory. In the last few years, the application of the AdS/CFT correspondence to the Superconducting/Superfluid phase has been an interesting subject, which states that the Einstein-Maxwell-scalar model provides a holographically dual description of a superconductor [2–4]. For reviews on holographic superconductors, see [5–7].

Recently, in Ref. [8], the authors proposed a gravity model which provides a holographically dual description of a 2+1 dimensional Josephson Junction. By choosing spatial dependence μ and non-vanishing constant J , which correspond to chemical potential and superfluid currents on the AdS boundary [9–12], respectively, the authors have solved numerically the coupled, nonlinear partial differential equations and reproduced the well-known result that the DC Josephson current is proportional to the sine of the phase difference across the junction. Meanwhile, they also studied dependence of the maximum current on the temperature and size of the junction, which matches precisely with the results for condensed matter physics.

In this paper, we would like to extend the work [8] to a holographic description for the (3+1)-dimensional holographic DC Josephson junction. We will study the 5D Einstein-Maxwell-scalar model and solve numerically the coupled, nonlinear partial differential equations of motion, and look for how to construct a holographic Josephson junctions by choosing the boundary conditions.

The paper is organized as follows. In Sec. 2, we review 5D Einstein-Maxwell-scalar theories and set up a gravity dual of a (3+1)-dimensional SNS Josephson junction. In Sec. 3, we show numerical results of the EOMs and study the characteristics of the (3+1)-dimensional holographic Josephson junctions. The last section is devoted to conclusion.

2 The model

In this section, we study the dynamics of a complex scalar field coupled to a U(1) vector potential in (4+1)-dimensional AdS spacetimes. The action of gravity can be written as

$$S = \int d^5x \sqrt{-g} \left(R + \frac{12}{L^2} \right), \quad (2.1)$$

where L is the curvature radius of asymptotic AdS Spacetimes. We are interested in the solution of the planar Schwarzschild black hole with the form:

$$ds^2 = -f(r)dt^2 + \frac{dr^2}{f(r)} + r^2(dx^2 + dy^2 + dz^2), \quad f(r) = \frac{r^2}{L^2} - \frac{M}{r^2}, \quad (2.2)$$

where M is the mass of the black hole. The black hole has the event horizon r_H at

$$r_H = M^{\frac{1}{4}} L^{\frac{1}{2}}. \quad (2.3)$$

The Hawking temperature of the black hole reads as

$$T = \frac{1}{4\pi} \left. \frac{df}{dr} \right|_{r=r_H} = \frac{M^{1/4}}{\pi L^{3/2}}. \quad (2.4)$$

In this background, we now consider a Maxwell field and a charged complex scalar field, with the action

$$S = \int d^5x \sqrt{-g} \left[-\frac{1}{4} F^{\mu\nu} F_{\mu\nu} - |\nabla\psi - iA\psi|^2 - m^2|\psi|^2 \right], \quad (2.5)$$

where $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$. The equations of motion include the scalar equation

$$-(\nabla_a - iA_a)(\nabla^a - iA^a)\psi + m^2\psi = 0, \quad (2.6)$$

and Maxwell's equations

$$\nabla^a F_{ab} = i[\psi^*(\nabla_b - iA_b)\psi - \psi(\nabla_b + iA_b)\psi^*]. \quad (2.7)$$

Taking a static ansatz

$$\psi = |\psi|e^{i\varphi}, \quad A = A_t dt + A_r dr + A_x dx, \quad (2.8)$$

where $|\psi|$, φ , A_t , A_r , and A_x are real functions of r and x , one can introduce the gauge-invariant field $M = A - d\varphi$, which is the same definition as that in [8].

With the black hole background (2.2) and the above ansatz (2.8), the equation of scalar field (2.6) can be written as the real part:

$$\partial_r^2 |\psi| + \frac{1}{r^2 f} \partial_x^2 |\psi| + \left(\frac{f'}{f} + \frac{3}{r} \right) \partial_r |\psi| + \frac{1}{f} \left(\frac{M_t^2}{f} - f M_r^2 - \frac{M_x^2}{r^2} - m^2 \right) |\psi| = 0, \quad (2.9)$$

and the imaginary part:

$$\partial_r M_r + \frac{1}{r^2 f} \partial_x M_x + \frac{2}{|\psi|} \left(M_r \partial_r |\psi| + \frac{M_x}{r^2 f} \partial_x |\psi| \right) + \left(\frac{f'}{f} + \frac{3}{r} \right) M_r = 0. \quad (2.10)$$

The imaginary part of the scalar equation can also come from the conservation of the source in Maxwell's equations. Meanwhile, Maxwell's equations (2.7) can be written as

$$\partial_r^2 M_t + \frac{1}{r^2 f} \partial_x^2 M_t + \frac{3}{r} \partial_r M_t - \frac{2|\psi|^2}{f} M_t = 0, \quad (2.12)$$

$$\partial_x^2 M_r - \partial_r \partial_x M_x - 2r^2 |\psi|^2 M_r = 0, \quad (2.13)$$

$$\partial_r^2 M_x - \partial_r \partial_x M_r + \left(\frac{f'}{f} + \frac{1}{r}\right) (\partial_r M_x - \partial_x M_r) - \frac{2|\psi|^2}{f} M_x = 0, \quad (2.14)$$

where a prime denotes derivative with respect to r . Because Eqs. (2.9)-(2.14) are coupled nonlinear equations, one can not solve these equations analytically. However, it is straightforward to solve them numerically.

In the following part of this paper, we choose the mass of the scalar field to be $m^2 = -3$, which is above the Breitenlöhner-Freedman bound [14]. In order to solve these coupled equations, first, we need to impose regularity at the horizon and the boundary on the radial coordinate. At the horizon $r = r_H$, $M_t(r_H) = 0$; At the Ads boundary, the scalar field takes the asymptotic form

$$|\psi| = \frac{\psi^{(1)}(x)}{r} + \frac{\psi^{(2)}(x)}{r^3} + \mathcal{O}\left(\frac{1}{r^4}\right), \quad (2.15)$$

the asymptotic behavior of the Maxwell fields are

$$M_t = \mu(x) - \frac{\rho(x)}{r^2} + \mathcal{O}\left(\frac{1}{r^3}\right), \quad (2.16)$$

$$M_r = \mathcal{O}\left(\frac{1}{r^3}\right), \quad (2.17)$$

$$M_x = \nu(x) + \frac{J}{r^2} + \mathcal{O}\left(\frac{1}{r^3}\right), \quad (2.18)$$

here, μ , ρ , ν and J are interpreted as the chemical potential, charge density, superfluid velocity and current, respectively [13, 15]. J is a constant. For $|\psi|$, both $\psi^{(1)}$ and $\psi^{(2)}$ are normalizable, one can impose the condition either $\psi^{(1)}$ or $\psi^{(2)}$ vanishes. For simplicity, we will consider the case $\psi^{(1)} = 0$ and interpret $\langle \mathcal{O} \rangle = \psi^{(2)}$ as the condensate. Here, we introduce the same phase difference $\gamma = \Delta\varphi - \int A_x$ as the case in [8], which is the gauge invariant and can be rewritten as

$$\gamma = - \int_{-\infty}^{\infty} dx [\nu(x) - \nu(\pm\infty)]. \quad (2.19)$$

Second, on the spatial coordinate x , we impose the Dirichlet-like boundary condition on M_r and Neumann-like boundary condition on $|\psi|$, M_t and M_x at $x = 0$. At $x = \pm\infty$, the function is x -independent. Thus the boundary conditions of the coupled equations (2.9)-(2.14) are determined by J and μ .

According to the study of paper [8], we can also introduce the critical temperature T_c of the junction, which is proportional to $\mu_\infty = \mu(\infty) = \mu(-\infty)$:

$$T_c = \frac{1}{\pi} \frac{\mu(\infty)}{\mu_c}, \quad (2.20)$$

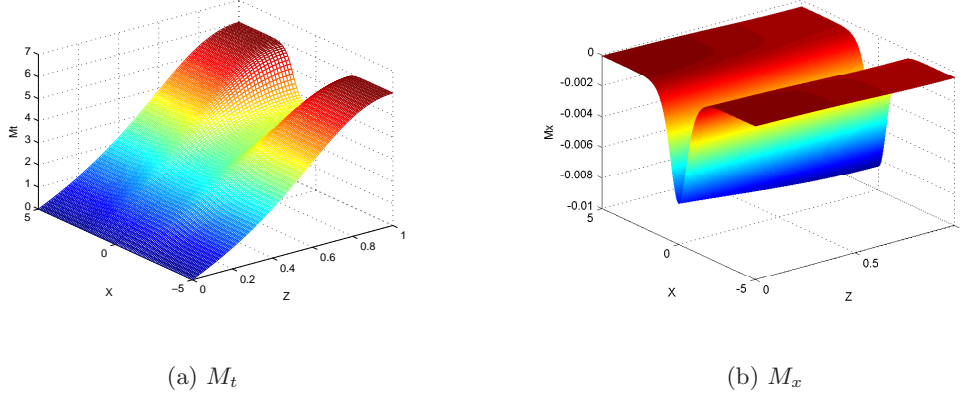


Figure 1. Maxwell fields M_t and M_x . The parameters are set to $J/T_c^3 = 0.0103$, $\mu_\infty = 6$, $L = 3$, $\epsilon = 0.6$, $\sigma = 0.5$.

where $\mu_c \approx 4.16$. Furthermore, the effective critical temperature inside the gap is

$$T_0 = \frac{1}{\pi} \frac{\mu(0)}{\mu_c} . \quad (2.21)$$

In order to describe an SNS Josephson junction, we choose the below $\mu(x)$ which is the same as that in [8]:

$$\mu(x) = \mu_\infty \left\{ 1 - \frac{1 - \epsilon}{2 \tanh(\frac{L}{2\sigma})} \left[\tanh\left(\frac{x + \frac{L}{2}}{\sigma}\right) - \tanh\left(\frac{x - \frac{L}{2}}{\sigma}\right) \right] \right\} , \quad (2.22)$$

where μ_∞ is the chemical potential at $x = \pm\infty$, and L is the width of our junction. The quantities σ and ϵ control the steepness and depth of our profile, respectively.

3 Numerical results

In this section, we will computer the solution of the coupled equations (2.9)-(2.14) numerically. It is convenient to set the change of variables $z = 1 - r_H/r$ and $\tilde{x} = \tanh(\frac{x}{4\sigma})$. First, we show the solutions for M_t and M_x in Fig. 1, with $\mu_\infty = 6$, $L = 3$, $\epsilon = 0.6$, and $\sigma = 0.5$.

Then, in Fig. 2, we find that the superfluid current is proportional to the sine of the phase difference across the junction, namely, the red dots come from the numerical calculations match precisely with the the black solid sine curve. Analyzing the graph, we can obtain the maximum current across the junction:

$$J_{\max}/T_c^3 \approx 0.871. \quad (3.1)$$

In Fig. 3(a), the dependence of J_{\max} on the width of the gap is shown. The graph predicts an exponential decay with the growing width of the gap in J_{\max} :

$$J_{\max}/T_c^3 = A_0 e^{-\frac{\ell}{\xi}}. \quad (3.2)$$

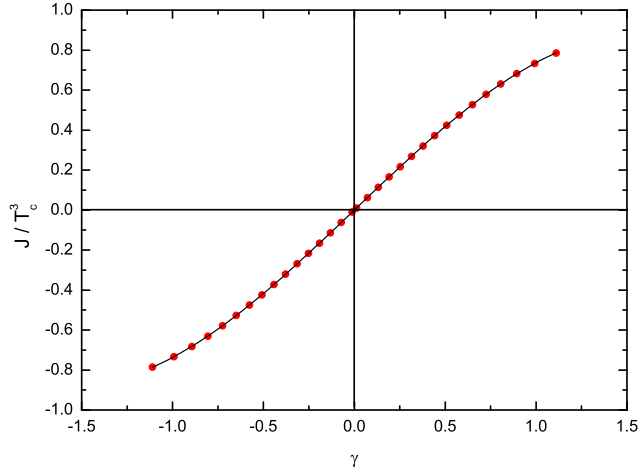


Figure 2. Superfluid current J_{\max}/T_c^3 as the function of the phase difference γ . The black line is the sine curve. The parameters are set to $\mu_\infty = 6$, $L = 3$, $\epsilon = 0.6$, $\sigma = 0.5$.

In Fig. 3(b), the condensate $\langle \mathcal{O} \rangle = \psi^{(2)}$ at zero current is shown. The graph also predicts an exponential decay with the growing width of the gap in J_{\max} :

$$\langle \mathcal{O} \rangle_{x=0, J=0}/T_c^3 = A_1 e^{-\frac{\ell}{2\xi}}. \quad (3.3)$$

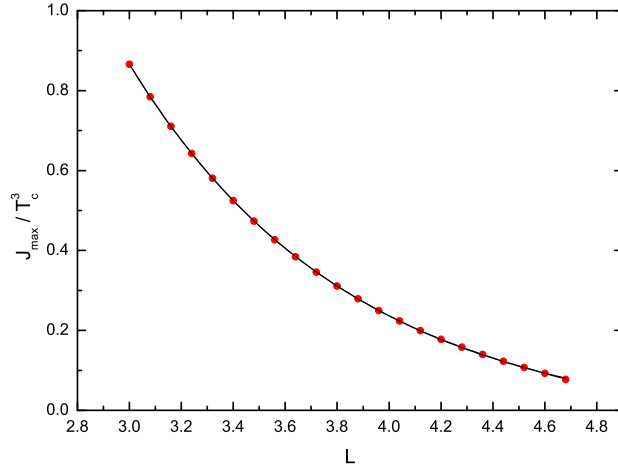
Fitting Eq. (3.2) and Eq. (3.3) with the two sets of data, we can obtain $\{\xi, A_0\} \approx \{0.85, 30.32\}$ and $\{\xi, A_1\} \approx \{0.84, 104.33\}$ for Eq. (3.2) and Eq. (3.3), respectively. The disagreement of two values of ξ is smaller than the 4D case in [8].

In Fig. 4, we obtain the relation of J_{\max} and T . Since $\epsilon = 0.6$, we show the region corresponds to $T/T_c < 0.6$, which depicts the character of an SNS Josephson Junction.

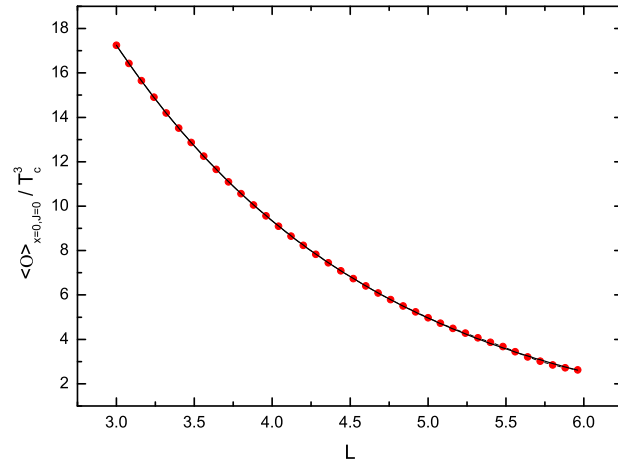
4 Conclusion

In this paper, the holographic description for the holographic DC Josephson junction is extended to the case of the (3+1)-dimensional junction. In the (4+1)-dimensional gravity background, we investigate a set of five couple, partial differential equations. By choosing spatial dependence μ and non-vanishing constant J , we find the the superfluid current is proportional to the sine of the phase difference across the junction. Moreover, the graph which predicts an exponential decay with the growing width of gap in J_{\max} is obtained. At last, the relation of J_{\max} and T is also obtained, from which one can verify that, near the critical temperature T_c , J_{\max} can reach zero.

From the result of this paper, one can see that the model of the holographic DC Josephson junction is valid for the (3+1)-dimensional case. It would be interesting to investigate the holographic DC Josephson junction in diverse dimensions and in various gravity theories.



(a)



(b)

Figure 3. The dependence of J_{\max} and $\langle O \rangle_{x=0}$ on L . The parameters are set to $\mu_{\infty} = 6$, $\epsilon = 0.6$, $\sigma = 0.5$.

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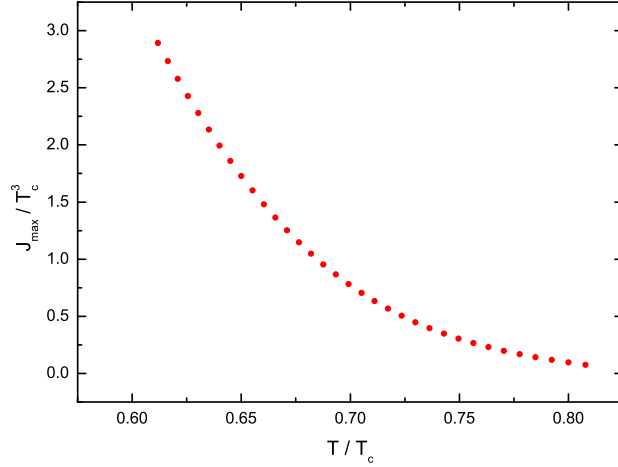


Figure 4. The dependence of J_{\max} on T , The parameters are set to $\epsilon = 0.6$, $\sigma = 0.5$

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